

## RESEARCH ARTICLE

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## Key Points:

- More than 1% perturbed PKP<sub>BC</sub>-PKP<sub>DF</sub> residuals are observed in each hemisphere
- We verify lateral variations of isotropic V<sub>p</sub> on regional scales in the UIC
- Possible effects from outer core and lowermost mantle structures are considered

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## Regionally heterogeneous uppermost inner core observed with Hi-net array

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**Abstract** Studies of velocity structure in the Earth's inner core and its volumetric variations illuminate our understanding of inner core dynamics and composition. Here we use an extensive number of seismograms recorded by the Hi-net array to construct complete empirical curves of PKP<sub>BC</sub>-PKP<sub>DF</sub> differential traveltimes. The nature of these curves implies that significant variations in traveltimes are accumulated during the passage of PKP<sub>DF</sub> waves through the uppermost inner core and rules out outer core structure effects. Uniform cylindrical anisotropy of a plausible strength in the uppermost inner core can also be ruled out as a cause of the observed traveltime variations because the range of sampled ray angles is too narrow. The configuration and strength of inhomogeneities from a recent tomographic model of the lowermost mantle cannot account for the observed traveltime variations. Therefore, we infer that either variations of *P* wave isotropic velocity on the scale of about hundred km and less are present in the uppermost inner core or the material may be organized in distinctive anisotropic domains, and both of these features may be superimposed on long-wavelength hemispherical structure. If the former holds true, the absolute magnitude of required *P* wave velocity perturbations from referent values is  $0.60 \pm 0.10\%$  in the quasi-eastern and  $1.55 \pm 0.15\%$  in the quasi-western hemisphere ( $0.85 \pm 0.05\%$  and  $1.10 \pm 0.10\%$ , respectively, with the lowermost mantle correction). The existence of these variations is a plausible physical outcome given that vigorous compositional convection in the outer core and variations in heat exchange across the inner core boundary may control the process of solidification.

### 1. Introduction

Since the discovery of the inner core (IC) [Lehmann, 1936], various spherically symmetric (1-D) *P* wave velocity models of the IC have been proposed [e.g., Jeffreys, 1939; Müller, 1973; Qamar, 1973; Anderson and Hart, 1976; Dziewonski and Anderson, 1981; Kennett and Engdahl, 1991; Kennett et al., 1995]. Although these models present an accurate description of a wide range of geophysical data, observational evidence from a number of studies supports a laterally heterogeneous IC. Various scales of IC heterogeneities have been reported in a better sampled, top part of the IC: quasi-hemispheric [e.g., Tanaka and Hamaguchi, 1997; Niu and Wen, 2001; Yu et al., 2005], regional [e.g., Kaneshima, 1996; Stroujkova and Cormier, 2004; Krasnoshchekov et al., 2005], and fine-scale structures [e.g., Creager, 1997; Cormier and Li, 2002; Leyton and Koper, 2007a; Peng et al., 2008]. The hypothesis that the top of the IC is a mushy zone seems to be well established [e.g., Cao and Romanowicz, 2004], with the estimated grain size on the order of 1–2 km [e.g., Vidale and Earle, 2000] and the scale length of volumetric heterogeneities in the uppermost inner core (UIC) on the order of 1–10 km [Leyton and Koper, 2007a]. Even smaller grain size of several hundred meters was suggested by Calvet and Margerin [2008] and Cormier and Li [2002] estimated several hundred meters near the top of the IC, increasing with depth up to tens of km. Both of these estimates come from interpreting inner core *Q* as a scattering *Q* that broadens the PKP<sub>DF</sub> pulse by transferring scattered high frequency into the later coda. Deguen et al. [2007] suggest that due to uncertainties in conjunction with iron conditions at high temperature/pressure, it is difficult to estimate the thickness of a mushy zone.

Apart from the process of solidification, the IC is experiencing a postsolidification deformation due to an internal stress field [e.g., Jeanloz and Wenk, 1988; Yoshida et al., 1996] and it is not entirely clear which one of these two processes has a dominant impact on crystal growth and IC texture. Cormier [2007] and Cormier et al. [2011] argue that an outward-stretched (radial) or inner core boundary (ICB) parallel-stretched (tangential) heterogeneity in the UIC can reconcile some observations of IC hemisphericity in seismic attenuation and scattering from PKiKP wave coda [Leyton and Koper, 2007b]. Their results are in line with the hypothesis that the quasi-eastern hemisphere (QEH) is freezing faster than the quasi-western hemisphere (QWH) [Aubert et al., 2008].

Although in odds with this scenario, a hypothesis of *Monnerieu et al.* [2010] and *Alboussière et al.* [2010] that the QEH is melting faster than the QWH is still consistent with the observed attenuation dichotomy between the two hemispheres of the IC [e.g., *Cao and Romanowicz*, 2004]. Crystal size in the QEH might be larger than typical wavelengths of body waves sensitive to IC structure, and this can explain observed lack of strong backscattering [*Leyton and Koper*, 2007b], while the melting component can explain higher attenuation. An alternative model would be an isotropic distribution of scale lengths in places where PKiKP coda has been observed in the mid-Pacific, and possibly a larger crystal (or organized patch of crystals) in regions where no strong PKiKP coda is observed [*Cormier et al.*, 2011].

That the above-described dichotomy in seismic velocity (and seismic attenuation) exists between the two hemispheres of the IC is largely based on the traveltimes and amplitudes of body waves sampling the IC in the equatorial region. In a recent study, however, *Ohtaki et al.* [2012] demonstrate that the IC does not have a simple hemispherical variation. Their analysis, while not global but based on unique raypath sampling of the IC under Antarctica, supports a single eyeball-shaped anomaly concentrated onto a small region beneath eastern Asia.

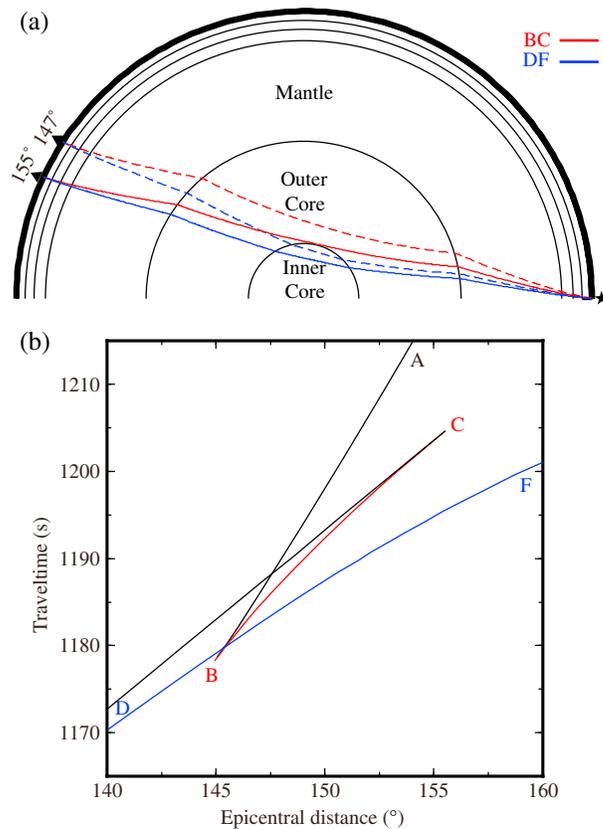
Much less constrained by seismic data than seismic velocity and attenuation is seismic anisotropy. *Poupinet et al.* [1983] observed that PKP<sub>DF</sub> waves travel faster along the quasi-polar than the quasi-equatorial paths through the IC. A concept of a uniform cylindrical IC anisotropy stems from this observation and was used to explain traveltime anomaly observed in PKP<sub>DF</sub> waves [*Morelli et al.*, 1986] and anomalous splitting of normal modes sensitive to IC structure [*Woodhouse et al.*, 1986]. Since then various seismological studies have been conducted in conjunction with IC anisotropy on a global scale [e.g., *Shearer and Toy*, 1991; *Tromp*, 1993; *Romanowicz et al.*, 1996; *Song and Helmberger*, 1998], its quasi-hemisphericity [e.g., *Tanaka and Hamaguchi*, 1997; *Creager*, 1999; *Garcia and Souriau*, 2000; *Niu and Wen*, 2002; *Irving and Deuss*, 2011], its radial variation [e.g., *Ishii and Dziewonski*, 2002; *Calvet et al.*, 2006; *Cormier and Stroujkova*, 2005] and its regional variation [e.g., *Isse and Nakanishi*, 2002; *Helffrich et al.*, 2002; *Tkalčić*, 2010]. For more details, see recent review papers [*Souriau*, 2007; *Tkalčić and Kennett*, 2008; *Cormier et al.*, 2011; *Deguen*, 2012].

A recent progress in normal mode techniques considering mode coupling [e.g., *Deuss et al.*, 2010] is in agreement with findings from body waves that point to complex structure in the IC superimposed on simpler long-wavelength average structure. Normal modes, however, are sensitive to relatively long spatial scales, comparable with the IC radius and hence constrain average structure, much less insights in fine structure that might characterize the UIC. Therefore, body wave techniques continue to be an essential tool in our quest to better understand geodynamical processes near the boundary between the solid and the liquid part of the core.

In contrast to normal modes, seismological techniques associated with body wave traveltime data have remained relatively unchanged over the last several decades. For example, an advantage of measuring differential traveltimes over absolute traveltimes has been demonstrated in numerous studies, and this remains one of the most exploited techniques since the pioneering work [e.g., *Cormier and Choy*, 1986]. The computational power is still insufficient to synthesize the seismic wavefield at 1 Hz and above for inner core-related seismic phases including the full range of 3-D influences in the mantle and inner core.

Thus, while it appears that seismological studies still have a long way to go before the reconciliation of various results can be made, the characterization and mapping of heterogeneity in the IC remains an important topic. Fortunately, what has changed dramatically in the last decade is the number of seismic recorders around the world, and therefore the capacity and potential to measure more traveltimes of core-sensitive body waves. Moreover, high dense arrays and array-based techniques probably represent future of seismic studies of the deep Earth's interior. For a more detailed account on array studies, see a review by *Rost and Thomas* [2002]. This work is entirely based on a new data set collected by a dense array of seismic stations with a goal of mapping structure in the UIC. It gives a glimpse of what can be achieved in deep Earth seismology with a large number of newly collected data and some of the well-established techniques. It leads to a conclusion that the UIC is regionally heterogeneous.

Our array of choice with unprecedented density of elements, all installed in boreholes, is the Hi-net array of short period instruments located in Japan. The array has significantly higher signal-to-noise ratio in comparison with the stations at the surface. We exploit the position of the Hi-net array and the corresponding regions of the UIC sampled by the raypaths of PKP waves recorded with the Hi-net array elements (section 2). We collect PKP traveltimes corresponding to raypaths that span epicentral distance range from 147° to 155° such that both PKiKP (PKP<sub>DF</sub>) and PKP<sub>BC</sub> phases are clearly recorded on the same seismogram (Figure 1). The finally selected data



**Figure 1.** (a) Raypaths of PKP<sub>DF</sub> (blue lines) and PKP<sub>BC</sub> (red lines) at epicentral distance of 147° (dashed lines) and 155° (solid lines). Star and triangle symbols indicate epicenter and station locations. (b) Traveltime curves of PKP waves around a triplication predicted from ak135 model [Kennett *et al.*, 1995].

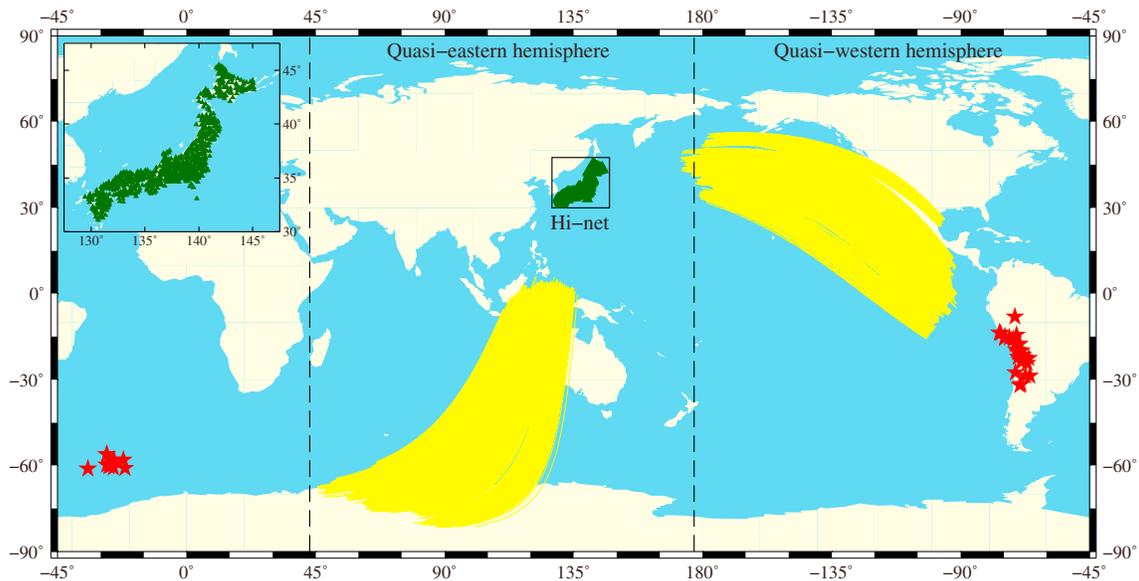
set consists of 47 individual events (from initial 242 ones), whose raypaths provide a dense coverage of the IUC. This allows us to construct complete empirical traveltime curves of PKP<sub>BC</sub>-PKP<sub>DF</sub> differential traveltimes (hereafter called BC-DF). We then analyze BC-DF residuals with respect to spherically symmetric reference models of the Earth (section 3).

First, we exclude a possibility that uniform cylindrical anisotropy in the IUC has a significant effect on the observed BC-DF traveltimes. This follows from the fact that the span of the angles between the raypaths of PKP<sub>DF</sub> in our data set and the rotation axis of the Earth is too narrow for anisotropy to have any noticeable effect (section 4.1). Our analysis results in 47 individual 1-D structural models (velocity profiles as a function of depth) that represent distinct volumes of the IUC. When these profiles are considered together, they are in agreement with a hemispherical dichotomy in the IUC (section 4.2), but also argue for regional variations in isotropic velocity with around 1% amplitude in each quasi-hemisphere (section 4.3). We verify by forward synthetic tests that different velocity profiles in the lowermost outer core (section 4.4), heterogeneities in the

outer core and the lowermost mantle (section 4.5), which could all hypothetically have significant impact on BC-DF, do not explain the observed variations. We then discuss the observed velocity profiles and their lateral variation in the context of previous seismological and geodynamical work (section 5) and conclude with ideas for future work directions (section 6).

## 2. Data

The data considered in this study are collected from vertical component seismograms recorded at the high-sensitivity seismograph network (Hi-net) consisting of about 800 stations operated by the National Research Institute for Earth Science and Disaster Prevention in Japan. The high-sensitivity seismometers are installed at depths of 100 m or deeper inside the boreholes to observe stable seismic waveforms by reducing surface ground noise. The Hi-net array elements cover the entire area of Japan with an average spacing of about 20 km [Okada *et al.*, 2004]. The configuration of Hi-net array and the location of large teleseismic earthquakes beneath southern Indian and northeastern Pacific Oceans deliver raypaths with near-equatorial sampling of the IC that span the epicentral distance range between 147° and 155° and sample both QE and QW of the IC. The fact that the Hi-net array is superabundant and favorably positioned to allow the recording of the core-sensitive wavefield presents an unprecedented opportunity to study the Earth's core "with a high-resolution lens." Several studies focusing on anisotropy, attenuation, and other seismic structures of the IC using Hi-net data demonstrated the reliability and merit of various observations in the investigation of the deep Earth structure [e.g., Kawakatsu, 2006; Kazama *et al.*, 2008; Wookey and Helffrich, 2008].



**Figure 2.** Surface projections of PKP<sub>DF</sub> raypaths considered in this study. Total of 47 events passed strict quality selection criteria and their corresponding raypaths are shown in red and yellow. Hi-net array stations (also shown enlarged in inset) are displayed by green triangles. The quasi-hemispheric boundaries (43°E, 177°W) defined by Tanaka and Hamaguchi [1997] are represented with dashed lines.

We collect waveform data from 242 events that occurred in and around South America and the Southern Atlantic Ocean between 2004 and 2009. All events are located at the epicentral distances ranging between 147° and 155° from the Hi-net array stations (Figure 2). Total of 47 events (Table 1), out of which 12 events in the QEH (No. 1–12 in Table 1) and 35 events in the QWH (No. 13–47 in Table 1), were selected applying restrictive data quality criteria based on the correlation coefficient, signal-to-noise ratio (SNR), and visual inspection of waveforms. Event information was taken from the reviewed Bulletin of the International Seismological Centre (ISC), which is manually checked by ISC analysts and relocated. The targeted events have the body wave magnitudes ( $m_b$ ) between 5.0 and 6.7. All qualified seismograms that were selected for a further study satisfy the following conditions: (1) the correlation coefficient between BC and DF waveforms is larger than 0.6; (2) the SNR of the PKP<sub>DF</sub> is higher than 2; and (3) the PKP<sub>DF</sub> traveltime curves are clearly identified by eye when all waveforms for an event are aligned with respect to the detected PKP<sub>BC</sub> phases (Figure 3). If the range of the epicentral distance for the PKP<sub>DF</sub> traveltime curves is too short (under 2°) or include empty spaces (over 1.5°), although the PKP<sub>DF</sub> traveltime curves were clear, we ruled the event out. Prior to the waveform cross correlation, all seismograms were band-pass filtered between 0.3 and 2 Hz and 6070 seismograms were finally selected from more than 50,000 seismograms.

Raypaths in this study almost exclusively sample each quasi-hemisphere because not only the bottoming points of the PKP<sub>DF</sub> phase are located within each quasi-hemisphere defined by previous studies (see Table 2) but also the corresponding entry and exit points at the ICB are almost entirely confined within a single hemisphere for most of earlier studies [e.g., Tanaka and Hamaguchi, 1997; Creager, 1999; Garcia and Souriau, 2000; Niu and Wen, 2001; Garcia, 2002; Oreshin and Vinnik, 2004; Irving and Deuss, 2011] (Table 2). Moreover, the entry, bottoming, and exit points are closely clustered in each quasi-hemisphere. Because we focus on regional variation of velocity in the UIC, the entry and exit points should be approached critically in contrast to most previous studies concerning just the bottoming points. The lateral extent of IC area for selected raypaths is 31.3° for the QEH and 43.0° for the QWH. At the ICB depth, these ranges correspond to about 665 and 915 km for each hemisphere, respectively.

### 3. PKP<sub>BC</sub>-PKP<sub>DF</sub> Differential Traveltimes

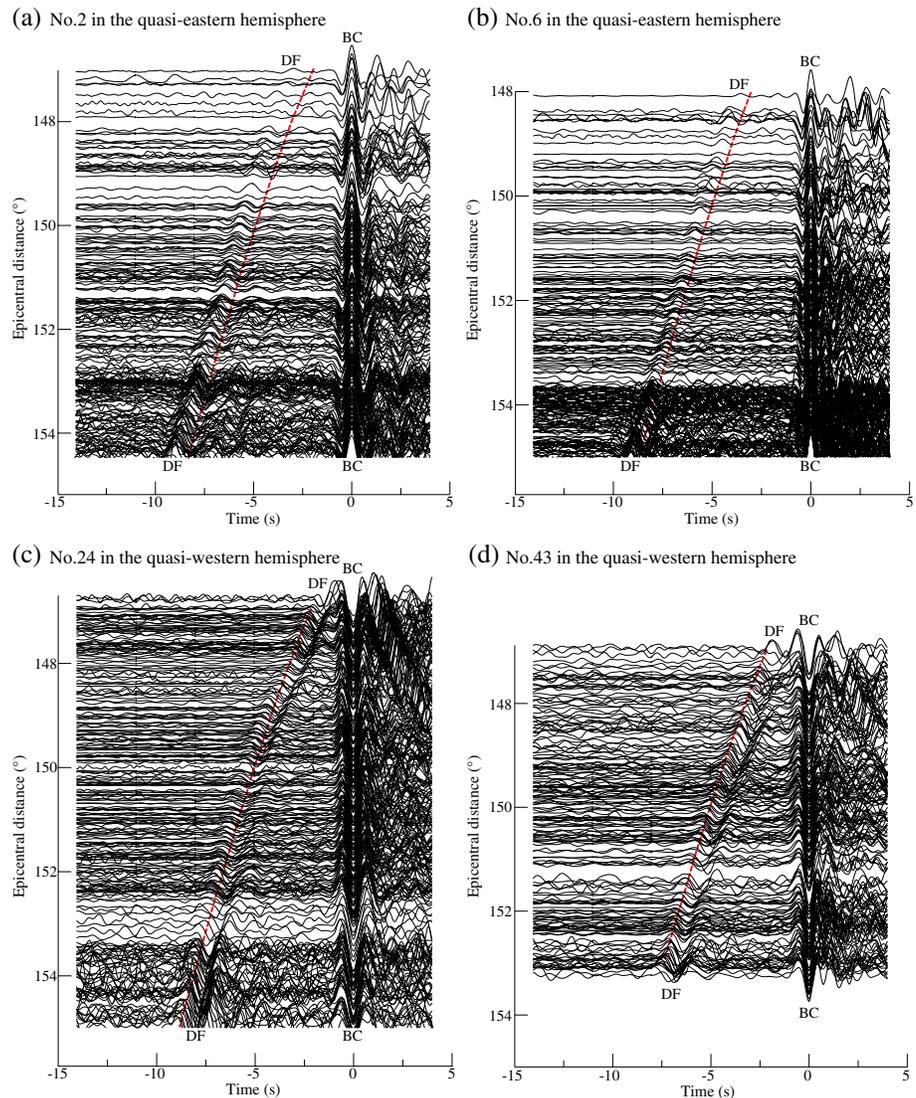
Since PKP<sub>DF</sub> and PKP<sub>BC</sub> phases have similar raypaths in the uppermost mantle (Figure 1a), PKP<sub>BC</sub>-PKP<sub>DF</sub> differential traveltimes (BC-DF) are only weakly sensitive to event depth and origin time uncertainties. We compared the reviewed ISC Bulletin and the Preliminary Determination of Epicenters Bulletin by the National Earthquake Information Center to estimate possible mislocation of earthquakes, and observed that the

**Table 1.** Event Parameters Used in This Study From the Reviewed Bulletin of the ISC

No	Origin Date <sup>a</sup>	Origin Time (UT)	$m_b$	Latitude, °N	Longitude, °E	Depth (km)
1	2006-01-02	06:10:48	6.2	-61.0106	-21.6489	10.0
2	2008-07-01	01:54:40	5.6	-58.2510	-21.8078	3.4
3	2004-09-11	21:52:34	5.9	-58.0207	-25.5207	35.6
4	2008-02-10	12:22:03	6.3	-60.7841	-25.5944	10.0
5	2005-08-04	12:11:18	5.7	-59.6285	-25.8459	44.3
6	2008-10-04	07:56:50	5.7	-59.4516	-25.8971	20.6
7	2006-05-29	05:20:38	5.5	-59.6523	-26.3024	37.4
8	2009-04-16	14:57:06	6.0	-60.2882	-27.0291	17.6
9	2007-07-31	02:42:49	5.8	-56.1028	-27.7318	113.3
10	2008-07-30	20:15:12	5.2	-59.7161	-27.8818	136.2
11	2005-09-09	19:55:18	5.7	-56.0067	-27.9467	125.0
12	2006-08-20	03:41:45	6.2	-61.0160	-34.2851	0.5
13	2004-09-07	11:53:04	6.0	-28.5709	-65.8815	22.1
14	2007-11-18	05:40:09	5.6	-22.6479	-66.2655	216.4
15	2005-06-02	10:56:00	5.5	-24.1334	-67.0103	194.9
16	2006-07-08	09:13:01	5.0	-28.6014	-67.4005	142.0
17	2005-11-17	19:26:54	5.9	-22.3676	-67.9421	161.7
18	2007-12-26	23:40:54	5.5	-22.3836	-68.3077	107.6
19	2006-11-14	22:22:12	5.4	-23.3092	-68.3195	109.5
20	2006-06-27	02:07:32	5.3	-22.5086	-68.4005	102.7
21	2004-06-17	01:16:00	5.2	-21.1833	-68.4451	115.4
22	2006-06-27	11:24:36	5.2	-21.2606	-68.4507	124.1
23	2008-01-05	07:29:33	5.5	-22.8861	-68.5037	105.4
24	2007-11-24	05:02:08	5.5	-23.7141	-68.8211	96.3
25	2005-08-14	02:39:39	5.6	-19.7422	-69.0809	114.6
26	2005-09-09	11:26:05	5.7	-31.6236	-69.1411	114.0
27	2008-09-10	16:11:59	5.5	-20.3049	-69.1785	3.9
28	2005-06-13	22:44:32	6.7	-19.9172	-69.2156	111.9
29	2009-05-08	13:44:53	5.6	-32.0330	-69.7032	113.3
30	2005-04-16	22:41:14	5.6	-17.6207	-69.7036	118.8
31	2007-11-14	15:40:49	6.6	-22.3208	-69.7803	33.6
32	2008-02-04	17:01:31	5.9	-20.1855	-69.9476	42.8
33	2007-03-24	19:13:50	5.6	-19.7094	-70.0611	27.7
34	2006-11-20	14:38:28	5.7	-17.6936	-70.1605	41.5
35	2005-07-13	12:06:11	5.5	-17.8565	-70.2123	79.3
36	2006-04-09	20:50:46	5.5	-20.4669	-70.2285	36.2
37	2009-04-17	02:08:09	5.8	-19.5806	-70.5024	28.7
38	2006-08-09	22:36:09	5.5	-14.4277	-70.6862	35.0
39	2004-08-27	00:43:52	5.7	-27.4305	-70.7903	30.6
40	2007-02-07	18:55:00	5.3	-17.7624	-70.9882	65.7
41	2007-07-21	13:27:03	6.1	-08.0859	-71.2071	633.7
42	2008-07-08	09:13:07	5.8	-15.9861	-71.7482	122.6
43	2005-07-26	14:11:34	5.7	-15.3386	-73.0597	109.1
44	2005-05-03	19:11:38	5.5	-14.8785	-74.6701	31.4
45	2007-06-28	08:44:48	5.4	-15.7255	-74.8709	27.9
46	2007-08-18	02:52:36	6.0	-13.8095	-76.2970	33.4
47	2007-08-19	01:22:38	5.8	-13.6101	-76.6104	13.4

<sup>a</sup>Dates are formatted as year/month/day.

maximum difference is 33.8 km in depth and 0.2° in epicentral distance. Uncertainties of the theoretical BC-DF traveltimes imposed by these maximum values are only 0.03 s taking into account the uncertainty in depth and 0.16 s taking into account the uncertainty in epicentral distance. In addition, the effects from strong heterogeneities near the core-mantle boundary (CMB) could be considered less significant under the assumption that the separation between  $PKP_{DF}$  and  $PKP_{BC}$  is smaller than a typical heterogeneity wavelength in the  $D''$  layer [e.g., Shearer and Toy, 1991; Bijwaard et al., 1998; Lay et al., 1998; Bréger et al., 1999; Romanowicz et al., 2003]. Indeed, a study of possible trade-off effects of the lowermost mantle structure and IC anisotropy on BC-DF confirmed that the trends observed in BC-DF residuals cannot be explained exclusively by the  $D''$  structure [Tkalčić et al., 2002].



**Figure 3.** Selected teleseismic waveforms recorded on Hi-net stations aligned with respect to the detected  $PKP_{BC}$  waves for (a) event “No. 2” and (b) event “No. 6” with ray sampling in the QEH and (c) event “No. 24” and (d) event “No. 43” with ray sampling in the QWH of the IC (see Table 1). Red dots are theoretical predictions of  $PKP_{DF}$  wave arrivals from ak135 model [Kennett *et al.*, 1995].

The waveform cross-correlation method was used to identify  $PKP_{BC}$  and  $PKP_{DF}$  phases and to measure BC-DF [e.g., Shearer and Toy, 1991; Tkalčić *et al.*, 2002]. The method takes advantage of the fact that the waveforms of those two core phases are similar because of a similar source radiation and the proximity of their paths through the Earth’s mantle. Before performing a cross correlation of  $PKP_{BC}$  and  $PKP_{DF}$  phases on the large number of Hi-net stations for a single event in an automatic fashion, we pick a clear  $PKP_{BC}$  phase manually. This single time then serves as a reference time. The seismogram for the reference pick corresponds to epicentral distance of about 149° to avoid interference with the  $PKP_{AB}$ ,  $PKP_{DF}$  and  $PKiKP$  phases that appear closer to  $PKP_{BC}$  phase when the epicentral distance approaches 145° on the one or 155° on the other side of the range (Figure 1b). We collect the  $PKP_{BC}$  phase arrivals from the rest of the seismograms by cross correlation with the reference  $PKP_{BC}$  phase. Another cross correlation is then performed to detect the  $PKP_{DF}$  phase on each seismogram. BC-DF are measured from cross-correlation lag times, the maxima of the cross-correlation function. Lastly, all selected seismograms are visually inspected. According to the reference model ak135 [Kennett *et al.*, 1995],  $PKP_{BC}$  and  $PKP_{DF}$  phases can be observed on the same seismogram for epicentral distances from about 145° to 155°. However, we do not consider data corresponding to the

**Table 2.** The Definition of Quasi-Hemispherical Boundaries From Previous Studies and Ray-Sampling Points From This Study

Quasi-Eastern Hemisphere	Quasi-Western Hemisphere	Reference
43°E–177°E	177°E–43°E	<i>Tanaka and Hamaguchi [1997]</i>
40°E–160°E	160°E–40°E	<i>Creager [1999]</i>
40°E–160°E	160°E–40°E	<i>Garcia and Souriau [2000]</i>
40°E–180°E	180°E–40°E	<i>Niu and Wen [2001]</i>
60°E–180°E	180°E–60°E	<i>Garcia [2002]</i>
50°E–120°W	120°W–50°E	<i>Oreshin and Vinnik [2004]</i>
14°E–151°W	151°W–14°W	<i>Irving and Deuss [2011]</i>
Range of bottoming points in the QEH	Range of bottoming points in the QWH	This study
95.70°E–130.54°E	114.14°W–138.01°W	
Range between entry and exit points	Range between entry and exit points	
45.40°E–135.53°E	171.99°E–91.15°W	

epicentral distances smaller than 147° because PKP<sub>BC</sub> and PKP<sub>DF</sub> arrivals are too close to each other, making the systematic detection of PKP<sub>DF</sub> phases by cross correlation less reliable.

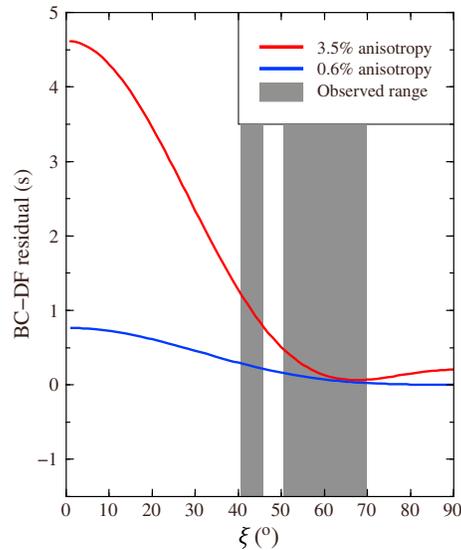
A traveltime residual  $\Delta t$  for each seismogram is defined as a difference between  $(BC - DF)_{obs}$ —the observed (measured) differential traveltime, and  $(BC - DF)_{1D}$ —the theoretical differential traveltime for different spherically symmetric reference Earth models. This can be described by the following equation:

$$\Delta t = (BC - DF)_{obs} - [(BC - DF)_{1D} + (BC - DF)_{ellip}] \quad (1)$$

where  $(BC - DF)_{ellip}$  is the Earth’s ellipticity correction (not significant, but still taken into account). If the lowermost mantle structure is corrected for, the above equation becomes

$$\Delta t = (BC - DF)_{obs} - [(BC - DF)_{1D} + (BC - DF)_{ellip} + (BC - DF)_{LM}] \quad (2)$$

where  $(BC - DF)_{LM}$  represents contribution from the lowermost mantle. The lowermost mantle contribution will be discussed in more detail in section 4.5.



**Figure 4.** Ranges of  $\zeta$  (the angle between PKP<sub>DF</sub> and the Earth’s rotation axis in the IC) for the waveforms from this study are represented as gray-shaded areas. The theoretical curves of traveltime residuals for a uniform cylindrical anisotropy of 0.6% [Shearer and Toy, 1991] and 3.5% [Creager, 1992] are plotted in blue and red lines, respectively.

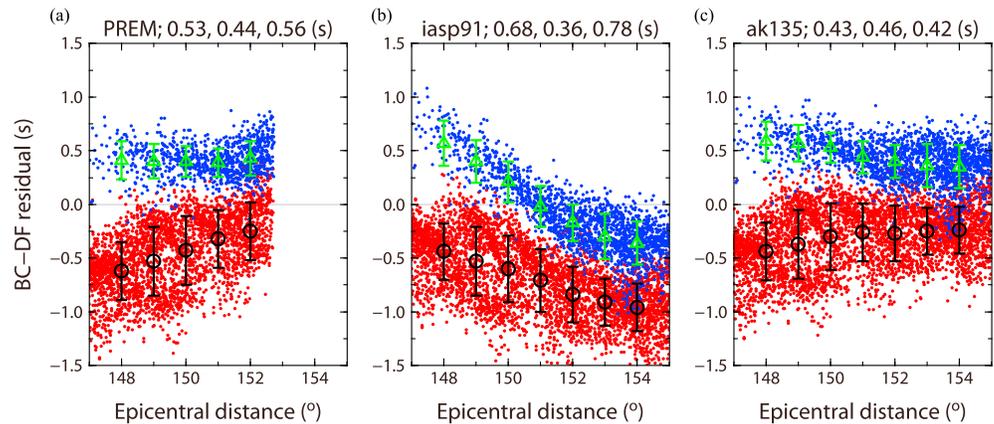
## 4. Results

### 4.1. Argument Against Uniform Cylindrical Anisotropy

PKP<sub>DF</sub> ray angles with respect to the Earth’s spin axis (conventionally denoted by  $\zeta$ ) for the raypaths in this study fall between 40.56° and 45.79° in the QEH and 50.42° and 69.83° in the QWH. The values of these angles indicate that all our paths are quasi-equatorial ( $\zeta$  between 35° and 90°). Assuming that the IC has a uniform cylindrical anisotropy about the Earth’s spin axis, the perturbed velocity can be expressed as

$$\delta Vp = A + B \cos^2 \zeta + C \cos^4 \zeta \quad (3)$$

For  $\zeta$  varying between 35° and 90°, the equation (3) predicts the range of BC-DF residuals to be about 1.7 s for 3.5% anisotropy ( $A = 0.018, B = -0.165, C = 0.557$  km/s) [Creager, 1992] and 0.4 s for 0.6% anisotropy ( $A = 0.000, B = 0.012, C = 0.056$  km/s) [Shearer and Toy, 1991] (see curves in Figure 4). However,  $\zeta$  corresponding to the selected waveforms in our study is restricted within narrow ranges of about 5° and 19° in each quasi-hemisphere (gray areas in Figure 4). It is straightforward to show that the effect of a possible strong uniform cylindrical anisotropy (such as that of 3.5%) is too small



**Figure 5.** Observed BC-DF residuals with respect to three spherically symmetric reference Earth models: (a) PREM [Dziewonski and Anderson, 1981], (b) iasp91 [Kennett and Engdahl, 1991], and (c) ak135 [Kennett et al., 1995] for the raypaths sampling in the QE (blue dots) and the QW (red dots). The means for the data confined within the interval  $\pm 1^\circ$  are calculated at each whole degree (green triangles for the QE and black circles for the QW). Corresponding  $\pm 1$  standard deviation is shown with vertical bars. RMS values are calculated for the entire data set, the QE data only, and the QW data only, and are shown by the three numbers above each plot.

to explain the magnitude of the observed BC-DF residuals. More specifically, for these small variations in angle  $\zeta$ , the corresponding predicted range of residuals for 3.5% anisotropy is less than 0.5 s in both hemispheres (see the extent of red lines in gray areas of Figure 4), while our observed traveltime residuals are much larger (Figure 5).

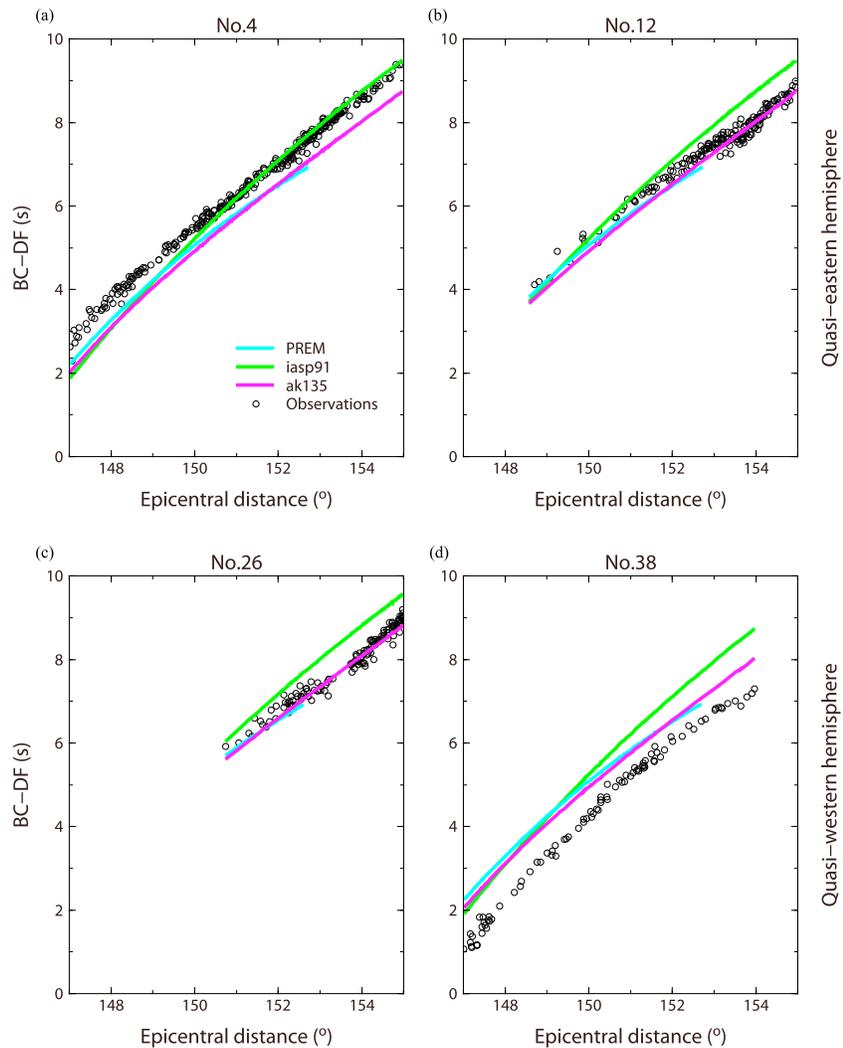
#### 4.2. Hemispherical Dichotomy in the UIC

When the BC-DF residuals for three commonly used reference models (PREM (preliminary reference Earth model) [Dziewonski and Anderson, 1981], iasp91 [Kennett and Engdahl, 1991], and ak135 [Kennett et al., 1995]) are plotted as a function of epicentral distance, these plots show clear distinction between the QE and the QW (Figure 5). Negative traveltime residuals are robustly observed in the QW, indicating a slower UIC with respect to the 1-D Earth reference models. Positive traveltime residuals in the QE indicate that the  $P$  wave velocity in the UIC is faster than the 1-D Earth reference models, and most observed residuals show positive values except for the iasp91 reference model. The residuals of the QE with respect to the iasp91 change from being positive to being negative as the epicentral distance increases because the slope of the observed and predicted BC-DF curves is quite different. This difference of the slope is mainly constrained by the velocity structure at the bottom of the outer core as will be shown by a synthetic test described in section 4.3. On the other hand, in case of PREM and ak135, the slope of the BC-DF curves is similar between the observation and prediction for each event (Figure 6).

The root-mean-square (RMS) values between the observations and the reference models are 0.53, 0.68, and 0.43 s for PREM, iasp91, and ak135, respectively. Since the slopes of the BC-DF curves between the observation and the prediction from iasp91 model show large deviations, we conclude that ak135 model is more relevant in further modeling of the velocity variations in the UIC. It should not be surprising that ak135 constrains PKP traveltimes after some adjustments better than iasp91, as ak135 was a significant improvement from iasp91 as a result of the addition of PKP traveltimes. The mean values and standard deviation of the observed BC-DF residuals are  $+0.41 \pm 0.16$  s for PREM and  $+0.42 \pm 0.19$  s for ak135 in the QE and  $-0.46 \pm 0.33$  s for PREM and  $-0.31 \pm 0.28$  s for ak135 in the QW. These mean values are quite comparable to the results obtained in previous studies [e.g., Tanaka and Hamaguchi, 1997; Niu and Wen, 2001].

#### 4.3. Lateral Variations on Regional Scales in the UIC

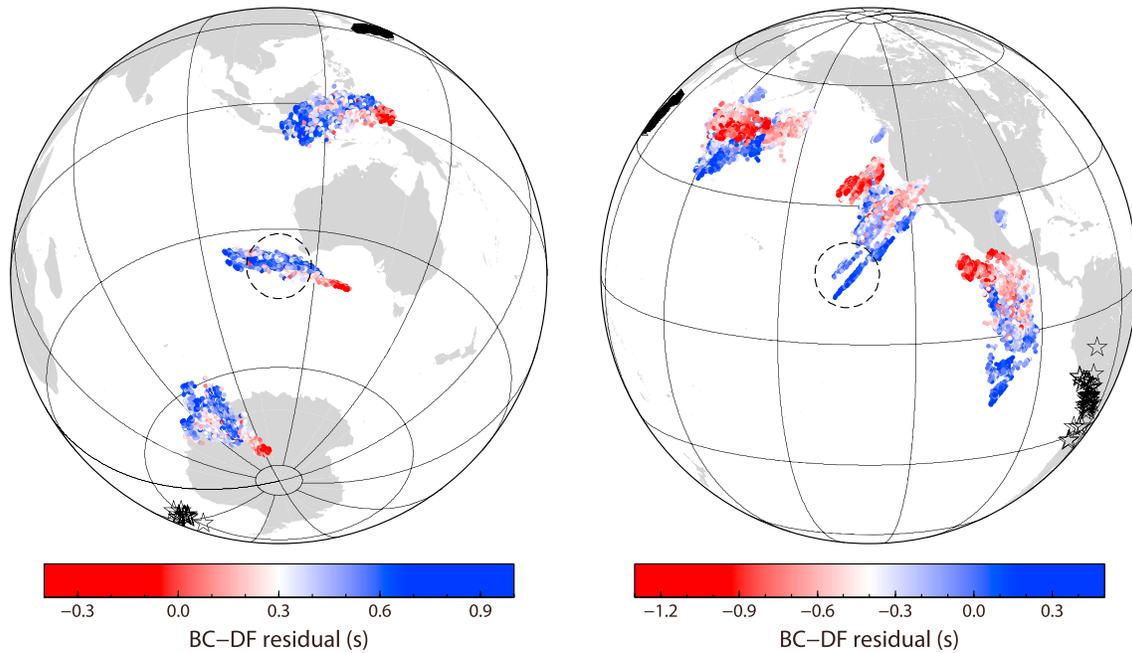
The shift of empirical BC-DF curves for all events (Figure 6) suggests that continuous variations of the velocity structure exist on regional scales. These variations could be confirmed at a glance from the distribution of residuals on the entry, bottoming, and exit points in the IC (Figure 7). All residuals are spread from  $-0.33$  s to  $+1.00$  s in the QE and from  $-1.25$  s to  $+0.46$  s in the QW. We identify a smooth transition from a slower eastern region to a faster western region in the QE. Similarly, there is an apparent transition from a slower northern region to a faster southern region in the QW. Traveltime residuals that are shown with



**Figure 6.** Theoretical BC-DF from PREM (cyan line), iasp91 (green line), and ak135 models (magenta line) are plotted against the observed BC-DF (black dots) for (a) the event No. 4 and (b) event No. 12 in the QEH, and (c) the event No. 26 and (d) event No. 38 in the QWH. The slopes of the PREM- and ak135-based BC-DF theoretical curves and observations are similar, but they differ for the iasp91 model.

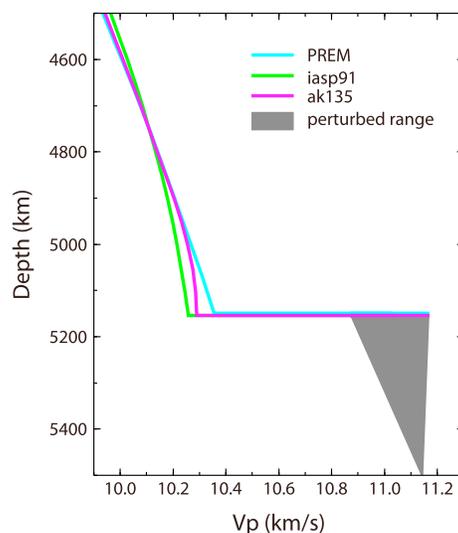
circles in Figure 7 are a result of an averaged velocity field within the corresponding ray's Fresnel zone, the value of which determines the PKP<sub>DF</sub> traveltime. Fresnel zones corresponding to adjacent sampling points cover similar volumes of the UIC, but they also each include a slightly different volume causing a smooth spatial variation in the observed residuals. The maximum radius of the Fresnel zone for a PKP<sub>DF</sub> phase with a dominant period of 1 s could be up to about 6.9° (148 km at the ICB; see the dashed circle in Figure 7) at the epicentral distance of 150° [Calvet *et al.*, 2006]. A lateral extent of the UIC sampled by raypaths in this study exceeds this radius in both sampling areas (31.3° in the QEH and 43.0° in the QWH).

In order to explain the observed shift of BC-DF curves, we perturb IC velocity structure using a grid search approach, starting from spherically symmetric *P* wave velocity models. PREM [Dziewonski and Anderson, 1981], iasp91 [Kennett and Engdahl, 1991], and ak135 [Kennett *et al.*, 1995] are used to represent the mantle and the outer core, while *P* wave velocity at the top of the IC is sampled from 10.87 km/s to 11.17 km/s with an interval of 0.01 km/s (Figure 8). Velocity structure between the top of the IC and the depth of about 5500 km (350 km below the ICB) is linearly interpolated. The PKP<sub>DF</sub> phases considered here do not travel through the IC below the depths of about 350 km, which corresponds to the epicentral distances of 155°. Therefore, the velocity structure outside these boundaries does not affect the BC-DF.



**Figure 7.** Distribution of residuals on the entry, bottoming, and exit points in the IC is plotted in each quasi-hemisphere. Dashed circle indicates the radius of the Fresnel zone for a PKP<sub>DF</sub> phase with a dominant period of 1 s at the epicentral distance of 150°. Black stars and triangles represent the selected events and Hi-net stations, respectively.

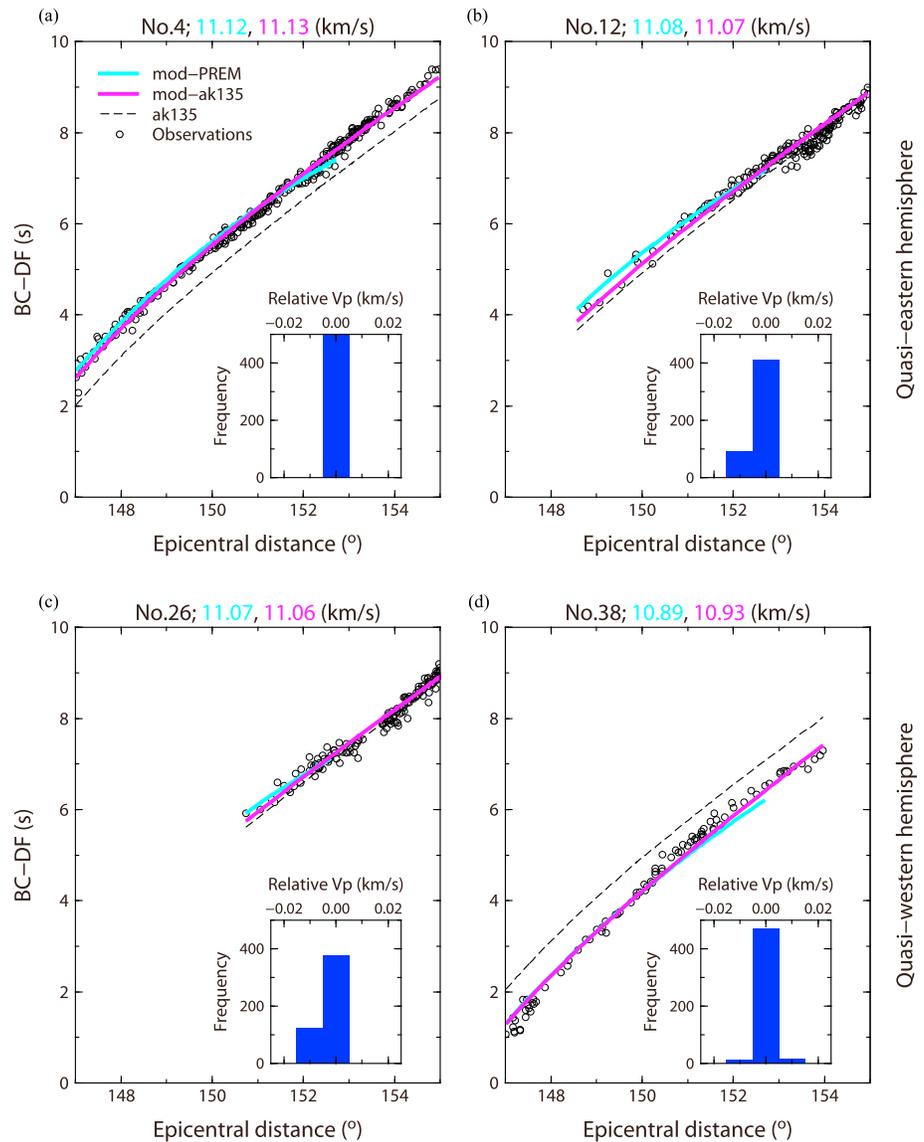
The best fitting models were determined by finding the lowest RMS value between the observed and theoretical BC-DF curves. The starting models are generated from each of the three reference models for each of the 47 selected events. Possible uncertainties due to sampling different event-station pairs within a broad region, event mislocations, site characteristics, different radiation patterns, and different source time functions, are negligible because we estimate a single *P* wave velocity model for each event. For simplicity, the best fitting modification of PREM for an individual event is referred to as mod-PREM and a consistent nomenclature is used for other two reference models (mod-ak135 and mod-iasp91). It is clear that the slopes



**Figure 8.** Trial models modified from PREM (cyan line), iasp91 (green line), and ak135 (magenta lines) for grid search. The velocity at the top of the IC was varied from 10.87 km/s to 11.17 km/s with an interval of 0.01 km/s (gray shadow).

of the best fitting BC-DF curves from the forward modeling do not change when different events are compared (Figures 9a and 9b for the QEH, and Figures 9c and 9d for the QWH the same as Figure 5 for comparison). The best fitting curves for mod-PREM and mod-ak135 show little difference and have similarly low RMS values. Using a bootstrap method, we resampled the observations and performed a grid search 500 times to estimate the robustness of the best fitting models. On average, about 78% of the best fitting models were reselected (insets in Figure 9).

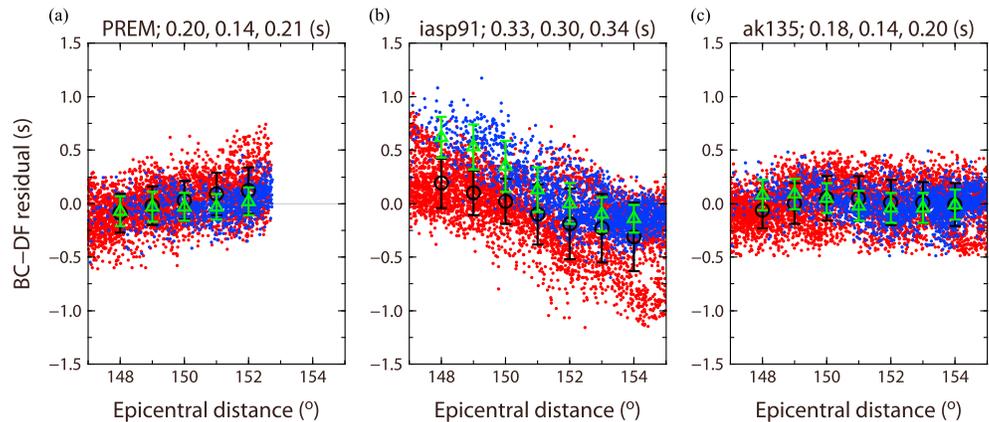
For all analyzed events, the RMS value decreased from 0.53 s for PREM to 0.20 s for mod-PREM and from 0.43 s for ak135 to 0.18 s for mod-ak135 (Figure 10) and the best fitting velocity for each event is specified in Table 3. Since each individual path is treated separately, it is not surprising that these values represent significant improvement in velocity structure when compared with the reference models. The RMS values for the QEH and the QWH are 0.14 s and 0.21 s for mod-PREM models and 0.14 s and 0.20 s for mod-ak135 models. At the same time, even the best fitting models derived from iasp91 do not explain



**Figure 9.** Observed (black dots) and theoretical BC-DF curves for ak135 (black dashed line), mod-PREM (cyan line), and mod-ak135 (magenta line) for (a) event No. 4, (b) event No. 12, (c) event No. 26, and (d) event No. 38 (see also Figure 6). The best fitting velocities at the ICB according to mod-PREM and mod-ak135 are noted above plots. Insets show the bootstrap results as histograms with the total number of models on vertical and the compressional velocity perturbations on horizontal axes.

the observations in contrast with other modified models. A decrease in the residual RMS value for the best fitting models modified from iasp91 is irrelevant because the BC-DF curves between the observations and the reference model are quite dissimilar as mentioned in section 4.1. Therefore, we conclude that individual mod-ak135 models and mod-PREM models can explain the observations equally well.

For individual mod-ak135 and mod-PREM models, the best fitting *P* wave velocities at the top of the IC vary between 11.07 and 11.13 km/s (about +0.3 to +0.8% with respect to ak135) and between 11.07 and 11.15 km/s (about +0.4 to +1.1% with respect to PREM) in the QEH. The velocities vary between 10.91 and 11.06 km/s (about -1.2 to +0.2% with respect to ak135) and between 10.87 and 11.07 km/s (about -1.4 to +0.3% with respect to PREM) in the QWH (Figure 11). In other words, the velocity in the UIC is perturbed in 2.0–2.5% on a global scale, considering ak135 and PREM. Moreover, on shorter spatial scales, the maximum velocity perturbation is about 0.5–0.7% in the QEH and 1.4–1.7% in the QWH even for an equatorial path with narrow



**Figure 10.** BC-DF residuals calculated with respect to the best fitting models: (a) mod-PREM, (b) mod-iasp91, and (c) mod-ak135 for the QEH (blue dots) and the QWH (red dots). The means for the data confined within the interval  $\pm 1^\circ$  are calculated at each whole degree (green triangles for the QEH and black circles for the QWH). Corresponding  $\pm 1$  standard deviation is shown with vertical bars. RMS values are calculated for the entire data set, the QEH data only, and the QWH data only and are shown by the three numbers above each plot.

range (Figure 12). These amplitudes are significant in comparison with the results from the previous studies, especially in the context of the proposed cylindrical anisotropy in the IC ranging from values below 1% to up to about 4%.

The best fitting velocity in Figure 12 represents the averaged velocity through a whole raypath from an entry to an exit point, not an independent velocity at the entry, bottoming, and exit points. Since not all Hi-net stations are aligned along the same great circle path from the event, PKP raypaths are not all in the same plane and hence they are not projected onto the same great circle. However, azimuthal variations of the stations are relatively small: from  $5.81^\circ$  to  $15.54^\circ$  (mean value is  $13.30^\circ$ ) for the events sampling the IC in the QEH and from  $4.54^\circ$  to  $18.42^\circ$  (mean value is  $10.87^\circ$ ) for the events sampling the IC in the QWH. We provide averaged 1-D velocity models since the number of station-event pairs is still insufficient to constrain more complex modeling approaches. Our study as it stands provides positive proof for the existence of regional velocity variation in the UIC. Velocity structure of the UIC is likely more complicated than that obtained by conventional approaches in which only bottoming points are considered using the angle  $\zeta$ . The structures at the ICB entry and exit points should have a significant effect on traveltimes of PKP<sub>DF</sub> waves.

#### 4.4. Possible Effects From Outer Core Structure

One possibility that has to be excluded before concluding that the bulk of traveltime anomalies can be attributed to the UIC structure is that the observed traveltime variations are affected by the structure in the outer core. We conduct several synthetic tests to the effects due to velocity perturbations in the outer and the inner core. The velocity in test models based on PREM is changed in the lowermost part of the outer core and the top of the IC. More precisely, the velocity structure between the depth of about 4470 km (680 km above the ICB) and the bottom of the outer core and the velocity structure from the top of the IC and the depth of about 5500 km (350 km below the ICB) were linearly interpolated (Figure 13a). On the one hand, the depth of 680 km above the ICB marks the upper bound of the depths sampled by the PKP<sub>BC</sub> phases in our study as it corresponds to the epicentral distances of  $145^\circ$ . On the other hand, the PKP<sub>DF</sub> phases considered here do not travel through the IC below the depths of about 350 km, which corresponds to the epicentral distances of  $155^\circ$ . Therefore, the velocity structure outside these boundaries does not affect the BC-DF.

The most remarkable result of synthetic tests is that when the velocity at the top of the IC is altered (solid and dashed red lines in Figure 13a) keeping the velocity in the outer core fixed to PREM values (dotted line), the predicted curves of BC-DF simply shift along the traveltime axis without significant alterations in the slope (Figure 13b). If, however, the velocity at the bottom of the outer core is altered (solid and dashed blue lines), while the velocity in the IC is kept the same (PREM; dotted line), the slope of the predicted curves of BC-DF changes, and the velocity gradient in the outer core affects the location of the c-cusp (Figure 13b). The

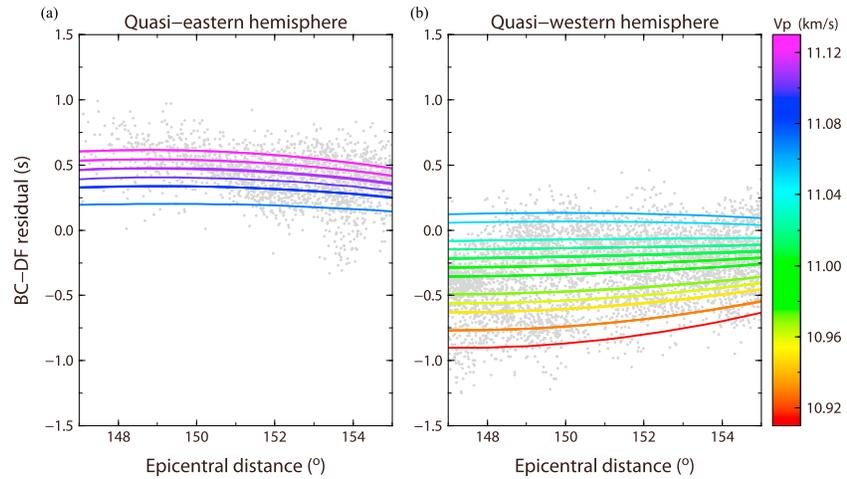
**Table 3.** Best Fitting Velocities From Forward Modeling

No	mod-PREM <sup>a</sup>		mod-ak135 <sup>a</sup>	
	$V_p^b$	$V_p$ for D'' corrected <sup>b</sup>	$V_p$	$V_p$ for D'' corrected
1	11.09	11.14	11.09	11.14
2	11.10	11.09	11.11	11.09
3	11.11	11.11	11.12	11.11
4	11.12	11.17	11.13	11.17
5	11.07	11.10	11.09	11.10
6	11.09	11.11	11.10	11.11
7	11.09	11.11	11.11	11.11
8	11.11	11.14	11.12	11.14
9	11.15	11.11	11.13	11.11
10	11.11	11.13	11.11	11.12
11	11.13	11.08	11.09	11.08
12	11.08	11.12	11.07	11.12
13	11.01	11.01	11.03	11.02
14	10.95	10.98	11.00	10.99
15	10.93	10.96	10.99	10.97
16	10.99	10.99	11.01	10.99
17	10.95	10.98	10.99	11.00
18	10.91	10.96	10.97	10.97
19	10.91	10.96	10.95	10.97
20	10.91	10.96	10.97	10.99
21	10.93	10.97	10.97	10.97
22	10.91	10.95	10.96	10.97
23	10.91	10.96	10.96	10.98
24	10.99	11.01	11.01	11.01
25	10.97	10.99	10.99	11.01
26	11.07	11.06	11.06	11.06
27	10.97	10.99	10.99	11.00
28	10.93	10.99	10.96	10.99
29	10.97	11.00	11.00	11.01
30	11.07	11.05	11.05	11.06
31	10.99	10.99	11.01	11.00
32	10.99	11.01	11.01	11.02
33	10.97	11.00	11.01	11.01
34	10.99	11.01	11.02	11.02
35	10.99	11.02	11.02	11.03
36	10.97	11.00	11.00	11.01
37	11.01	11.02	11.02	11.03
38	10.89	10.93	10.93	10.95
39	10.95	10.99	11.02	11.01
40	11.00	11.03	11.03	11.04
41	10.98	11.04	11.01	11.05
42	10.97	11.01	11.00	11.02
43	10.92	10.97	10.95	10.99
44	10.89	10.96	10.93	10.97
45	10.91	10.95	10.95	10.97
46	10.87	10.94	10.91	10.96
47	10.89	10.97	10.93	10.98

<sup>a</sup>The modification of PREM (ak135) best fitting the observed BC-DF for an individual event (same number as Table 1). The velocity structure is based on PREM (ak135) except between the top of the IC and the depth of about 5500 km (350 km below the ICB). For grid search,  $P$  wave velocity at the ICB was sampled from 10.87 km/s to 11.17 km/s with an interval of 0.01 km/s, and the velocity structure for 350 km below the ICB was linearly interpolated from the sampled  $V_p$ .

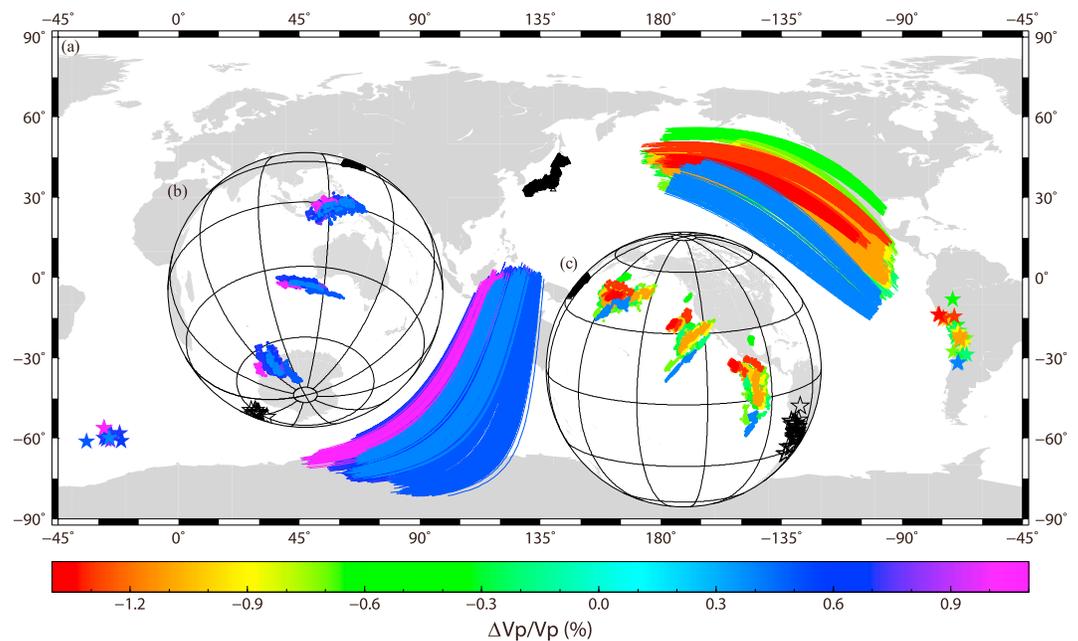
<sup>b</sup>The  $P$  wave velocity at the ICB from the best fitting model for the observed BC-DF and for the BC-DF considered an effect by the lowermost mantle anomalies.

observed traveltimes for all earthquakes selected in this study regardless of their location show similar slopes to that of ak135 and PREM, but their absolute values vary. This observation, according to our synthetic tests and assuming that reference Earth models provide reasonably good global average, indicates that outer core structure plays less important role in affecting traveltimes of PKP waves than structure in the UIC.

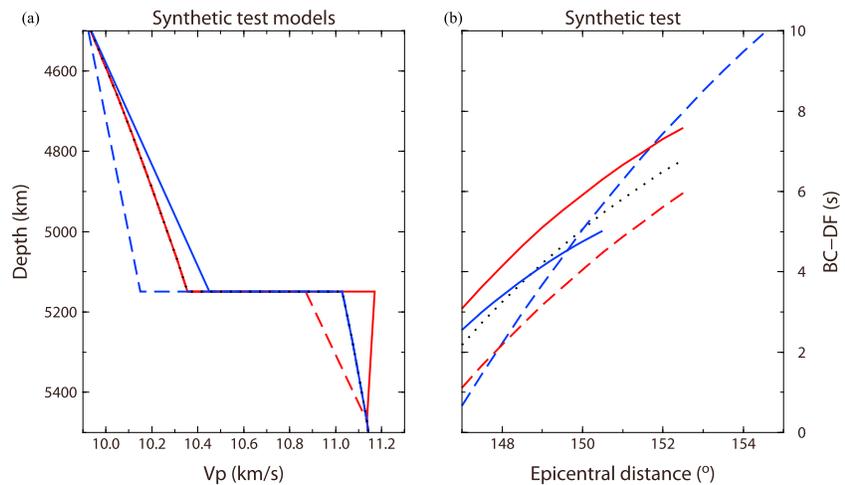


**Figure 11.** Variation of the theoretical BC-DF residuals for 47 events is shown in colors corresponding to best fitting velocities of mod-ak135 for (a) the QEH and (b) the QWH. Observed BC-DF residuals with respect to ak135 (same as in Figure 5c) are plotted with light gray dots.

Traditionally, the Earth’s outer core is considered well mixed (homogeneous relative to the passage of body waves) due to vigorous convection. Although several researches considered radial inhomogeneity at the top or the bottom of the outer core, it is still regarded laterally uniform [e.g., *Stevenson, 1987; Souriau and Poupinet, 1991; Romanowicz et al., 2003; Souriau et al., 2003; Tanaka, 2004, 2007; Ishii and Dziewonski, 2005; Zou et al., 2008*]. If velocity structure of the outer core is indeed uniform laterally, the observed difference in the empirical traveltimes considered here is best explained by lateral variations in IC velocity structure because the traveltimes of PKP<sub>BC</sub> and PKP<sub>DF</sub> raypaths are equally affected by outer core velocity structure and the slope of the empirical traveltimes curves stay unchanged.

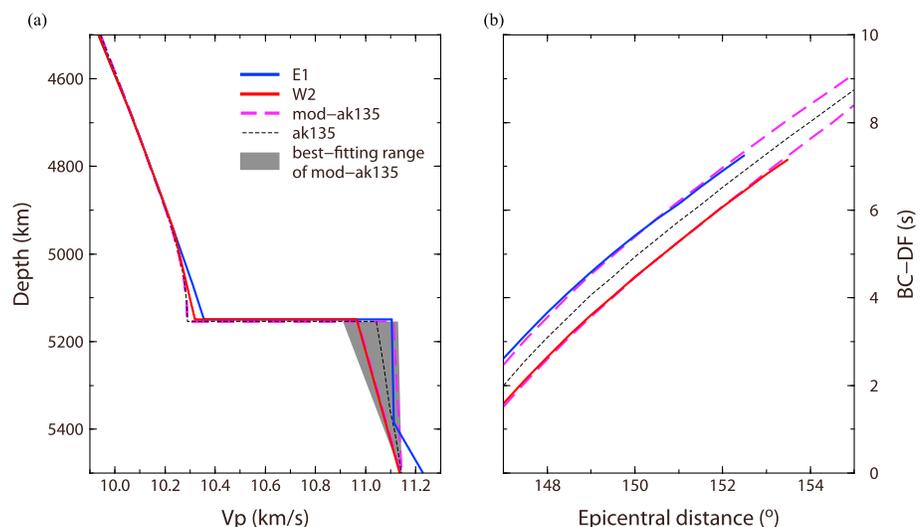


**Figure 12.** (a) Surface projections of great circle raypaths of PKP<sub>DF</sub> in the IC from this study. Colors represent the best fitting *P* wave velocity perturbation from PREM (defining mod-PREM). Stars and triangles represent the selected events and Hi-net stations, respectively. (b) The entry, bottoming, and exit points of PKP<sub>DF</sub> in the IC are plotted in the QEH using the same colors as in Figure 12a. (c) Same as Figure 12b but for the QWH.



**Figure 13.** Synthetic tests illustrating how changes in velocity profiles in the outer (blue line) and the inner core (red lines) affect theoretical BC-DF curves. (a) In the first test, the velocity profile in the outer core is changed (from a dashed to a solid blue line), while the profile in the IC is fixed to PREM values (dotted line). In the second test, the velocity profile in the IC is changed (from a dashed to a solid red line), while the profile in the outer core is fixed to PREM values (dotted line). (b) The BC-DF curves resulting from the profiles shown in Figure 13a with same color and line style.

It should be noted that *Yu et al.* [2005] suggested two different *P* wave velocity models at the bottom of the outer core for each quasi-hemisphere. According to their study, the QEH model, E1, has a PREM-like velocity structure while the QWH model, W2, has a lower velocity than PREM and a bit higher than ak135 in the bottom of the outer core (Figure 14a). Figure 14b shows 1 s difference between the E1 (blue) and W2 (red) models, which we confirm is mostly due to UIC structure. Namely, if we adhere to the ak135 velocity profile in the outer core, but change it at the ICB according to W2 and E1 models (two magenta lines), it becomes clear that BC-DF are not substantially affected by this dichotomy (Figure 14b). The difference in velocity structure between PREM and ak135 models at the bottom of the outer core (Figure 9) or among the models shown in Figure 14 does not yield significant difference in



**Figure 14.** (a) Different velocity profiles for the QEH and the QWH: ak135 (dashed black line), E1 (blue line), W2 (red line) [Yu and Wen, 2006], and two profiles of mod-ak135 (dashed magenta lines) showing similar velocity at the ICB with respect to E1 and W2. The range of best fitting mod-ak135 models for the events analyzed in this study is shaded with gray. The velocities at the ICB range from 10.91 to 11.13 km/s. (b) Theoretical BC-DF curves from the models shown in Figure 14a with same color and line style.

the slopes of the BC-DF curves. Other plausible velocity variations (such that other data set fits are not significantly reduced) at the bottom of the outer core would only have a minimal effect. Therefore, it appears most likely that regional variations of  $P$  wave velocity in the UIC have the most dominant effect on the observed empirical curves of BC-DF.

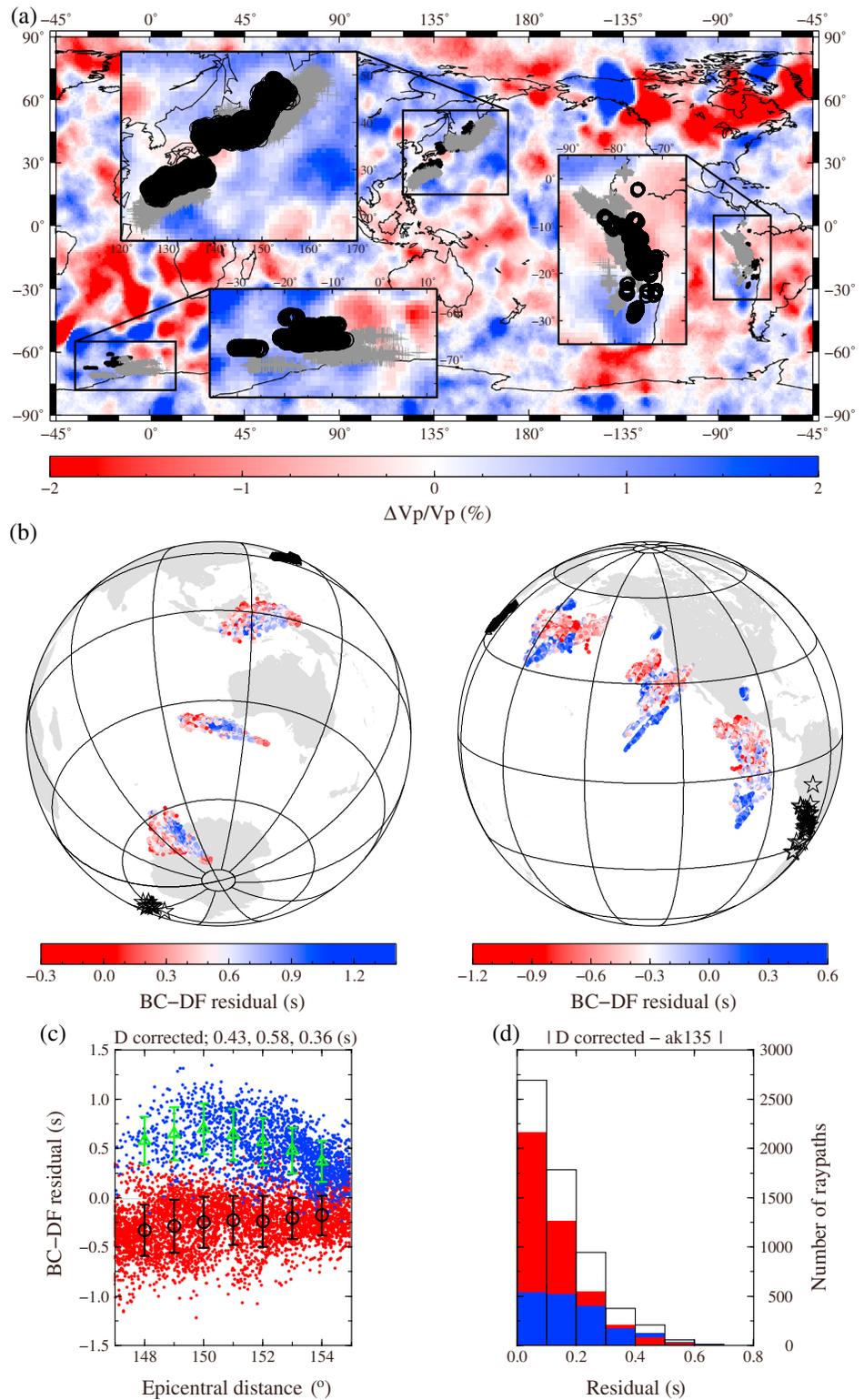
#### 4.5. Possible Effects From Mantle Structure

Most previous studies make an assumption that  $PKP_{DF}$  and  $PKP_{BC}$  waves transit the lowermost mantle in a way that their differential traveltimes are not significantly affected by heterogeneous structure. This assumption is valid in the upper mantle where the raypaths are less separated. However, various scales of heterogeneities in the lowermost mantle have been reported [e.g., Haddon and Cleary, 1974; Bataille and Flatte, 1988; Wysession, 1996; Su and Dziewonski, 1997; Garnero, 2000; Dziewonski et al., 2010]. These anomalies can influence traveltimes of PKP phases [e.g., Bréger et al., 2000; Tkalčić et al., 2002; Ishii et al., 2002; Romanowicz et al., 2003; Tkalčić, 2010]. The distance between piercing points of  $PKP_{BC}$  and  $PKP_{DF}$  waves at the CMB varies from about 263 km ( $4.3^\circ$ ) to about 474 km ( $7.8^\circ$ ) for the epicentral distance range  $147^\circ$ – $155^\circ$  considered in this study (Figure 15a). Thus, it is possible and likely that the two PKP phases will pass through quite different velocity field.

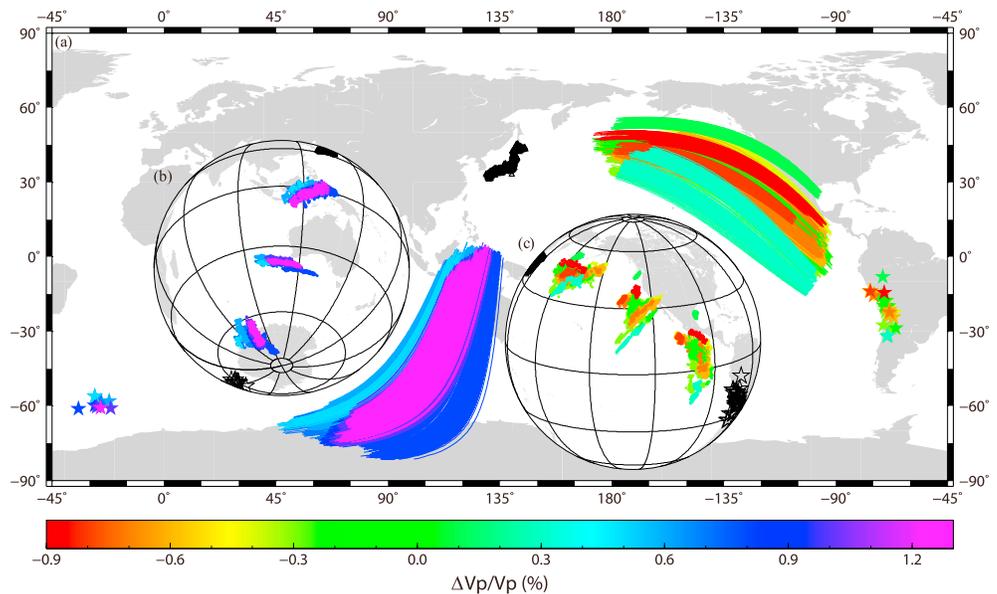
Different tomographic models disagree most strongly in the lowermost mantle, and the disagreements increase as the reported resolvable scale length decreases. Furthermore, the power heterogeneity spectrum of  $P$  wave velocity in the lowermost mantle estimated from previous tomography studies [e.g., Obayashi and Fukao, 1997; Bijwaard et al., 1998; Káráson and van der Hilst, 2001; Tkalčić et al., 2002; Li et al., 2008; Dziewonski et al., 2010; Soldati et al., 2012] is typical such that the  $P$  wave velocity variations are under  $\pm 1\%$ . There are rare exceptions, such as the study of Garcia et al. [2009] according to which the RMS velocity perturbations in the  $D''$  layer are 1.2%. A subjective choice of regularization is a possible cause for differences among traveltimes tomography models and the underestimate of the heterogeneity power. A recent tomography study of Young et al. [2013], in which the inversion is void of explicit regularization, shows a significant increase in the heterogeneity spectrum strength and a visible shift toward shorter spatial scales of heterogeneity in some areas of the lowermost mantle. This tomography model was derived using an extensive data set of different PKP raypaths from those that are utilized here, and, importantly, it is characterized by a good spatial coverage in the areas of the lowermost mantle critical for our consideration here, which is confirmed through resolution tests. Furthermore, the tomography model was critically constrained by the addition of PcP-P traveltimes that are sensitive to lowermost mantle structure. Therefore, in order to examine possible contribution to BC-DF by lowermost mantle structure ( $D''$  structure), the model of Young et al. [2013] was used to correct our data (Figure 15a).

We first recalculate the BC-DF residuals correcting for structure from the Young et al. [2013] tomography model. The distribution of residuals on the entry, bottoming, and exit points is plotted in Figures 15b and 15c in the same way as it was done in Figure 7. The BC-DF residuals show a spread from  $-0.23$  s to  $+1.35$  s in the QEH and  $-1.21$  s to  $+0.53$  s in the QWH (Figure 15b). The mean values and standard deviation of the  $D''$ -corrected residuals are  $+0.67 \pm 0.29$  s for PREM and  $+0.51 \pm 0.26$  s for ak135 in the QEH and  $-0.17 \pm 0.27$  s for PREM and  $-0.25 \pm 0.25$  s for ak135 in the QWH. The RMS values for the total, QEH, and QWH are 0.43, 0.58, and 0.36 s as against 0.43, 0.46, and 0.42 s for the case of using only ak135 model (Figures 15c and 5c). These quantitative changes show that the residuals are not only larger but also its range is broadened for the QEH, while the opposite is true for the QWH. However, a distinct presence of, and a smooth transition between slow and fast regions can still be discerned in each quasi-hemisphere.

Absolute values of the BC-DF difference between ak135 and  $D''$ -corrected model fall into the 0.0–0.3 s range for more than 89% of the paths (Figure 15d) which suggests that the lowermost mantle structure does not have a critical effect on BC-DF in this study. Theoretical maxima for BC-DF residuals with  $\pm 1$  and  $\pm 2\%$   $P$  wave velocity anomalies in the lowermost 300 km of the mantle are  $\pm 1.03$  s and  $\pm 2.06$  s. This means that the traveltimes of PKP phases could in theory be significantly affected by the lowermost mantle structure in certain regions. If such data are obtained from individual event-station pairs, it presents a challenge to separate the lowermost mantle from the IC effects. On the other hand, PKP data recorded on dense arrays and/or a group of adjacent events will lend themselves to much less ambiguous interpretation, still resting on an assumption that the separation between piercing points at the CMB is smaller than dominant length scales of anomalies in the lowermost mantle.



**Figure 15.** (a) The entry and exit points of PKP<sub>BC</sub> (gray crosses) and PKP<sub>DF</sub> (black circles) raypaths at the CMB plotted with the Young et al. [2013] tomography model in the background. (b) Distribution of residuals on the entry, bottoming, and exit points in the IC is plotted in each quasi-hemisphere. Black stars and triangles represent the selected events and Hi-net stations, respectively. (c)  $D''$ -corrected BC-DF residuals with respect to ak135, plotted in the same way as in Figure 5c. (d) Histogram of the absolute difference (in seconds) between the BC-DF traveltimes corrected for  $D''$  structure and traveltimes predicted by the ak135 model in the QEH (blue bar), QWH (red bar), and their global sum (white bar).



**Figure 16.** (a) Surface projections of great circle raypaths of  $PKP_{DF}$  in the IC from this study. Colors represent the  $D''$ -corrected best fitting  $P$  wave velocity perturbation from PREM (defining mod-PREM). Stars and triangles represent the selected events and Hi-net stations, respectively. (b) The entry, bottoming, and exit points of  $PKP_{DF}$  in the IC are plotted in the QE using the same colors as in Figure 16a. (c) Same as Figure 16b but for the QWH.

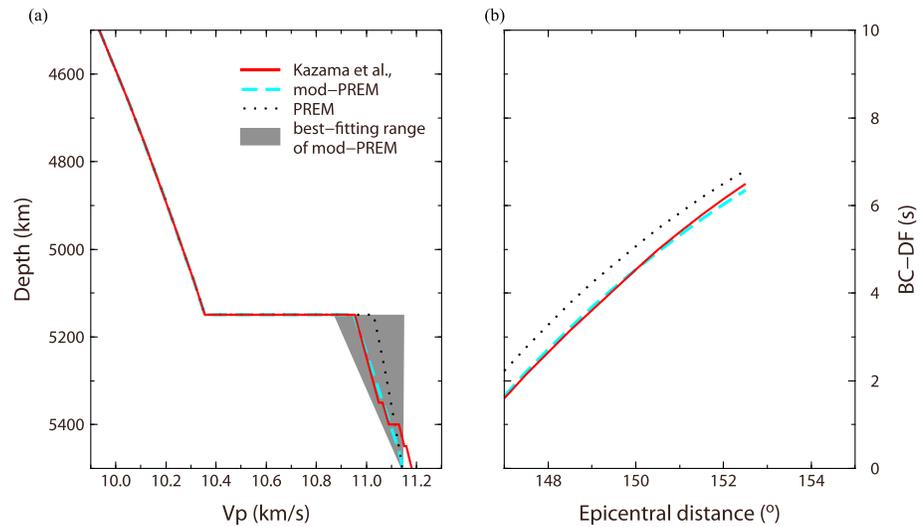
We perform forward modeling to quantitatively compare effects on BC-DF residuals due to  $D''$  structure for each event. The  $D''$ -corrected best fitting  $P$  wave velocities at the top of the IC vary between 11.08 and 11.17 km/s (about +0.3 to +1.2% with respect to ak135) and between 11.08 and 11.17 km/s (about +0.5 to +1.3% with respect to PREM) in the QE. The  $P$  wave velocities vary between 10.95 and 11.06 km/s (about  $-0.8\%$  to  $+0.2\%$  with respect to ak135) and between 10.93 and 11.06 km/s (about  $-0.9\%$  to  $+0.3\%$  with respect to PREM) in the QWH (Figure 16). The range of variation for the best fitting  $P$  wave velocity at the top of the IC increased in the QE and decreased in the QWH in comparison with the forward modeling that did not contain  $D''$  correction (Figure 12). However, an important observation is that the velocity of the UIC shows similar variation, whether we consider the  $D''$  effect or not (Table 3). The amplitude of variations in  $P$  wave velocity is 0.5–0.7% or 0.8–0.9% with  $D''$  correction in the QE and 1.4–1.7% or 1.0–1.2% with  $D''$  correction in the QWH. Correcting for lowermost mantle structure makes the difference between the two hemispheres less sharp but does not remove traveltime anomalies from our data set. Thus, although the heterogeneities in  $D''$  are an important contributor to traveltime variations, they cannot explain the observed significant variations of BC-DF in our study.

## 5. Discussion

### 5.1. Radial Velocity Structure in the UIC

Our modified models contain linearly interpolated velocity structure below the ICB similar to all three Earth reference models considered. However, some researches reported gradient transition in velocity or layered structures in the UIC [e.g., Yu and Wen, 2006; Kazama et al., 2008].

E1 and W2 [Yu and Wen, 2006] show different velocity structures not only in the bottom of the outer core but also in the UIC for each quasi-hemisphere (Figure 14a). Contrary to W2 that has a simple velocity structure in the IC (such as the case in our modified models), E1 has a sharp increase in the velocity gradient at 235 km beneath the ICB, which results in a changing slope of the BC-DF curves beyond the epicentral distances of  $151^\circ$ . However the BC-DF curves with respect to mod-PREM and mod-ak135 in the QE show little difference around the epicentral distance of  $151^\circ$  (Figure 14b). Considering more than  $\pm 0.5$  s deviation in the observed residuals, the changed gradients of the velocity structure are not significant.



**Figure 17.** (a) Different velocity profiles for the QWH: PREM (black dotted line), the velocity model of *Kazama et al.* [2008] (red line), and one profile of mod-PREM (dashed cyan line) showing similar velocity at the ICB with respect to the profile of *Kazama et al.* [2008]. The range of best fitting mod-PREM models for the events analyzed in this study is shaded with gray. The velocities at the ICB range from 10.87 to 11.15 km/s. (b) Theoretical BC-DF curves from the models shown in Figure 17a with the same color and line style.

Another BC-DF observation indicating regional complexity was also reported for the QWH by *Kazama et al.* [2008]. In contrast to a simple W2 model, their study showed more complex velocity structure for depth ranges 200–250 km, 250–300 km, and 300–350 km below the ICB. We also calculate the theoretical BC-DF from their model. In theory, a layered velocity structure causes discontinuous BC-DF curve, but their model does not produce significantly different effect than one of mod-PREM featured in Figure 17. We see no evidence for a discontinuous structure in the observed waveforms. In general, BC-DF curves cannot be explained by the complex model reported in *Kazama et al.* [2008], except for few events collocated with three events from their work.

The above difference in interpretation might stem from the fact that some of the existing models of IC radial structure have likely been obtained by gathering individual event-station pairs with different entry and exit points in the hemispheres or regions roughly defined by bottoming points of PKP<sub>DF</sub> wave. Some large residuals influencing a construction of alternative models could be caused by structure of the lowermost mantle or the existence of lateral variations in isotropic velocity in the UIC. However, considering that deviation of  $\pm 0.5$  s or more around the mean is fairly common in the observed BC-DF curves, many nonunique solutions are possible. There will be no significant difference between step-like (discontinuous) and radially smooth velocity models (interpolated linearly below the ICB using various Earth reference models). Since there is a lack of evidence in the waveforms for a step-like (discontinuous) structure, we argue that mod-ak135 and mod-PREM models present a plausible interpretation.

### 5.2. Lateral Variations in Isotropic Velocity in the UIC

Despite the fact that many studies have focused on its velocity structure, little is still known about regional variations of IC elastic properties. Part of the problem has been sparse raypath coverage and the lack of continuous spatial sampling on regional scales (several hundred km) within the IC. While it was possible to collect a data set of differential traveltimes of core-sensitive seismic phases over a range of epicentral distances from a large number of events [e.g., *Niu and Wen*, 2001; *Tkalčić et al.*, 2002; *Yu et al.*, 2005; *Garcia et al.*, 2006; *Waszek and Deuss*, 2011], it has not yet been possible to collect these measurements from a single event and use them in a similar fashion in which the traveltime data are used in reflection seismology. However, the establishment of dense permanent arrays with regional-scale apertures has opened a new era in the deep Earth seismology and some of its obvious pay-offs are seen in this study.

Sharp velocity gradients in lateral and radial direction in the UIC have been inferred before beneath Central America [e.g., *Creager, 1997*] and beneath the Indian Ocean [*Stroujkova and Cormier, 2004*]. Our results suggest that these regional variations in the isotropic velocity structure at the UIC are more widespread than what we have previously been able to infer. At this point, it is difficult to speculate if a larger magnitude variation inferred for the QWH has any geodynamical meaning because the number of sampling is different and its spatial distribution is restricted for both hemispheres. Although the difference in maximum amplitude of lateral perturbations of  $P$  wave velocity between the two hemispheres is reduced after the lowermost mantle corrections are applied, we note that these corrections are still not perfect. They come from a model which assumes that BC-DF data that participate in the model construction, though augmented with other types of data, are affected by exclusively mantle structure. Ideally, the lowermost mantle  $P$  wave velocity model would be constructed from mantle-sensitive data only; however at present, this is not possible as spatial coverage of the existing PcP-P waveform-correlated data is limited. It should also be noted that some extremely thin layers of large  $P$  wave velocity perturbations in the lowermost mantle have been documented through waveform modeling (ultralow velocity zones), but their thickness and spatial scales are beyond the resolution limit of tomographic imaging. Last but not least, the effects imposed by mantle structure will be fully understood only when they are quantitatively assessed with full waveform modeling. If it can be demonstrated from independent observations that the magnitude of velocity variation in the QWH is indeed stronger than in the QEH, then this, combined with robust observations from body waves and normal modes that the QWH is also slower, might provide important clues on melting and freezing processes. As of now, melting and freezing processes could both be invoked to explain velocity differences in both hemispheres [*Deguen et al., 2007*].

Since our PKP<sub>DF</sub> rays have mutually similar angles with respect to the spin axis of the Earth, significant variations of the BC-DF observed in our study cannot be explained with a model of uniform cylindrical anisotropy of the IC. Our data strongly suggest that small-scale regional variations (from the observed variation in empirical BC-DF curves) may be superimposed on a large-scale hemispherical pattern in isotropic velocity (from robust observations that one hemisphere is faster than another). Variations in convection and heat flux in the outer core are sources of heat transfer variations across the ICB, which influence IC growth and crystal alignment [e.g., *Yoshida et al., 1996*]. Regional variations in isotropic velocity in the IC have been suggested through the observation of variations in crystal alignment [*Creager, 1997*], texture [*Stroujkova and Cormier, 2004; Cormier, 2007*], and modeling of randomly oriented anisotropic patches [*Calvet and Margerin, 2008*]. A complexity in the observed traveltime data for polar PKP raypaths [e.g., *Leykam et al., 2010*] leads to a conclusion that the IC may be a conglomerate of anisotropic domains [*Tkalčić, 2010; Mattesini et al., 2013*] as an alternative possibility to describe a significant variation in traveltimes originating in the UIC.

## 6. Concluding Remarks

In this study, we have used recordings from about 800 stations of the Hi-net array stretching an elongated shape across the Japanese islands (over 15° in latitude) originated from teleseismic earthquakes whose PKP<sub>DF</sub> waves sample both “quasi-hemispheres” of the IC. This allowed us to construct individual complete empirical curves of PKP<sub>BC</sub>-PKP<sub>DF</sub> differential traveltimes (BC-DF) for a total of 47 events whose waveforms passed our strict quality selection criteria. These empirical BC-DF curves show many elements of unprecedented observations in modern seismology as they consist of traveltimes covering the theoretically predicted epicentral distance range, all from a single event. Through analysis of these traveltime curves, it becomes possible to critically examine and distinguish outer core and lowermost mantle effects from those that the IC structure imposes on traveltimes of PKP<sub>DF</sub> waves traversing its volume.

Our analysis reveals that both PREM and ak135 models could be adequate reference models for lowermost outer core velocity structure and that no significant modification of these models is necessary to explain the empirical BC-DF curves collected in this study. In addition, we have demonstrated that lowermost mantle structure is not responsible for variations in empirical BC-DF curves observed along the raypaths analyzed in this study. Instead, a likely solution that emerges as most effective to match the observations is to modify  $P$  wave velocity in the UIC along each raypath. Thus, our analysis leads to a conclusion that the IC could be

heterogeneous laterally on global (hemispherical pattern), regional (few hundred kilometers), and possibly local scales (several tens of kilometers). The latter estimate is close to the estimate of heterogeneity size published in studies of PKiKP scattering in the upper part of the IC where heterogeneity scale lengths were found to be on the order of 1 to 10 km [e.g., *Leyton and Koper, 2007a*]. We have quantified the minimum and maximum magnitude of lateral perturbations of *P* wave velocity with respect to ak135 and PREM to be of 0.5–0.7% (0.8–0.9% with the lowermost mantle correction) in the QEH and 1.4–1.7% (1.0–1.2% with the lowermost mantle correction) in the QWH.

The difference between the 1-D models reported in this array-based study and the existing spherically symmetric Earth models (as well as the models for the QEH and the QWH) implies that the imaging of the IC has reached a stage similar to the imaging of the mantle, where simple homogeneous structures no longer provide adequate prediction of the growing seismological data. A further refinement of existing models is thus necessary to represent lateral and radial variations on both regional and local scales. Future installations of dense arrays such as the Hi-net array and corresponding array studies will be a powerful tool to image and clearly distinguish between anisotropy- and isotropic velocity heterogeneity-related effects on traveltimes of core-sensitive body waves. This will then inevitably contribute to a better understanding of what the IC is made of and how it solidifies.

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